

# Solid methane toward deeply embedded protostars<sup>\*</sup>

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**Abstract.** We report on the detection of an absorption feature near  $7.67\ \mu\text{m}$  toward the deeply embedded protostellar objects W 33A and NGC 7538 : IRS9. Comparison with laboratory spectra shows that this feature can be identified as the  $\nu_4$  (‘deformation’) mode of solid state  $\text{CH}_4$  embedded in polar molecules (i.e.,  $\text{H}_2\text{O}$  and/or  $\text{CH}_3\text{OH}$ ) in icy grain mantles. The solid  $\text{CH}_4$  column density relative to solid  $\text{H}_2\text{O}$  is 0.4–1.9%. Unlike solid CO, solid  $\text{CH}_4$  does not seem to be subjected to out-gassing toward the warm source W 33A. The low gas-to-solid ratio for  $\text{CH}_4$  argues that  $\text{CH}_4$  is formed on grains either by hydrogenation of accreted C or by ultraviolet processing of  $\text{CH}_3\text{OH}$ -rich ices.

**Key words:** ISM: dust, extinction – ISM: molecules – ISM: abundances – infrared: ISM: lines and bands – stars: individual: NGC 7538: IRS9 – stars: individual: W 33A

## 1. Introduction

Theoretical models predict that molecular clouds may contain considerable quantities of methane ( $\text{CH}_4$ ) in the gas (e.g., Millar & Nejad 1985; Herbst & Leung 1986) or solid phase (e.g., Tielens & Hagen 1982). The extremely symmetric  $\text{CH}_4$  molecule can only be observed by its ro-vibrational transitions in the infrared, i.e. the  $\nu_3$  ‘stretching’ mode at  $3.32\ \mu\text{m}$  and the  $\nu_4$  ‘deformation’ mode at  $7.66\ \mu\text{m}$ . For solid  $\text{CH}_4$  these ro-vibrational lines merge into single broad bands, but somewhat shifted in wavelength due to interaction of the molecules in the matrix

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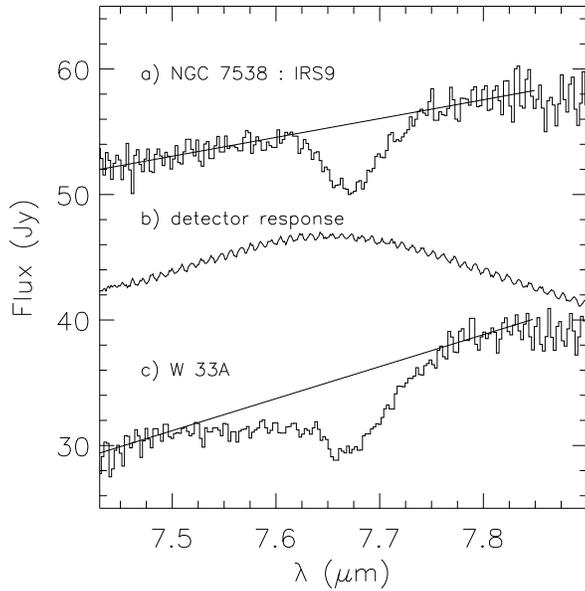
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and particle shape effects (Boogert et al. 1997, henceforth B97). Observations of both fundamental modes of interstellar  $\text{CH}_4$  are severely hindered by telluric absorption. Ground based and air-borne (KAO) observations of the  $7.66\ \mu\text{m}$  band have been done toward some bright protostellar sources. Lacy et al. (1991) detected solid and gaseous  $\text{CH}_4$ , and their detection toward W 33A was recently confirmed by KAO observations (B97). Both studies have to deal with rather poor signal-to-noise ratios due to telluric absorption or technical limitations.

With the Short Wavelength Spectrometer on board of the Infrared Space Observatory (ISO–SWS; Kessler et al. 1996; de Graauw et al. 1996) it is possible to observe solid and gaseous  $\text{CH}_4$  at high grating resolution ( $R \simeq 1800$ ), unhindered by the earth atmosphere. In this letter we present ISO–SWS spectra in the  $7.4\text{--}8.0\ \mu\text{m}$  range of the deeply embedded protostars W 33A and NGC 7538 : IRS9. For the first time we are able to compare absorption by interstellar solid  $\text{CH}_4$  in detail with laboratory studies of ices. The identification of the molecular environment of  $\text{CH}_4$  in grain mantles, its column density, as well as the gas-to-solid state ratio, will provide important clues to the formation history of the  $\text{CH}_4$  molecule and its importance for the chemical network in molecular clouds.

## 2. Observations

The protostellar objects W 33A and NGC 7538 : IRS9 were observed with ISO–SWS in full grating mode (‘AOT 6’). Both sources are much smaller than the ISO–SWS beam at  $7.7\ \mu\text{m}$  ( $14'' \times 20''$ ; Willner et al. 1982). Therefore, the resolution of the observations is comparable to the instrumental resolution for point sources at this wavelength, i.e.  $R \simeq 1800$ . The latest version of the standard ISO–SWS reduction procedure and calibration files was applied (OLP version 5.0 at 2/9/96; de Graauw et al. 1996; Schaeidt et al. 1996). The individual up and down detector scans of the ISO–SWS spectra show some systematic



**Fig. 1.** ISO-SWS spectra of the deeply embedded protostars NGC 7538 : IRS9 (a) and W 33A (c). The straight lines indicate the adopted continua. For comparison the detector response, scaled to the brightness of these sources, is shown (b).

differences in the flux level, as well as the spectral slope. This is probably caused by inaccuracies in the flux calibration, and the dark current subtraction (de Graauw et al. 1996). To minimize these systematic contributions to the noise, the detector scans were aligned by fitting straight lines and subsequently multiplying each scan to the mean level. This resulted in 10% lower RMS noise for W 33A and 70% for NGC 7538 : IRS9. Then the spectrum was resampled to the instrumental resolution with two points per resolution element. Figure 1 shows the final spectra. To check the reliability of the observed spectra, we have also plotted the detector response curve. No obvious residuals of the detector responsivity are apparent in the interstellar spectra.

### 3. Results

The spectrum of NGC 7538 : IRS9 shows a single, Gaussian shaped absorption feature near  $7.67 \mu\text{m}$ . The spectrum was converted to optical depth scale by adopting a local, straight line continuum, as indicated in Figure 1a. Then, the peak position, width (FWHM) and central optical depth were determined by fitting a Gaussian (Table 1). The spectrum of W 33A reveals a broad absorption between  $7.5\text{--}7.75 \mu\text{m}$ , that clearly is composed of two components (Fig. 1c). These distinct components were satisfactorily fitted by two Gaussians on optical depth scale (Table 1), using the continuum indicated in Figure 1c. We find that the narrow component of W 33A has the same peak position ( $7.675 \mu\text{m}$ ) and FWHM ( $0.062 \mu\text{m}$ ) as NGC 7538 : IRS9 within the fitting uncertainties. The broad feature in W 33A is centered on  $7.631 \mu\text{m}$ , and is evidently very weak or absent toward NGC 7538 : IRS9.

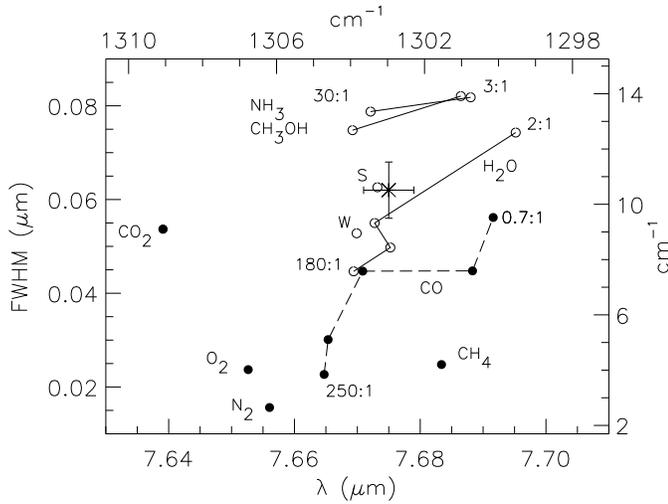
### 4. Identification of the observed features

An extensive laboratory study of interstellar solid  $\text{CH}_4$  in different ices and at different temperatures was made in B97. Their results are summarized in Figure 2, complemented with some laboratory data of Hudgins et al. (1993) for non-polar mixtures. It shows the peak position and width of the deformation mode of solid  $\text{CH}_4$  in a sample of ices after deposition at  $T=10 \text{ K}$ . The width and peak position are good discriminators between polar and non-polar matrices and the relative abundance of  $\text{CH}_4$  in the matrix. Upon warm-up to 30 and 50 K the separation of polar and non-polar ices in this diagram decreases, but is still present (B97).

A comparison of the ISO-SWS observations with the laboratory simulations shows that the observed  $7.67 \mu\text{m}$  feature is likely due to absorption by interstellar solid  $\text{CH}_4$ , embedded in polar molecules (Fig. 2). Especially  $\text{H}_2\text{O}:\text{CH}_4$  mixtures fit well. The solid line in Figure 2 shows that the laboratory profile broadens and shifts to larger wavelength with decreasing  $\text{H}_2\text{O}:\text{CH}_4$  abundance ratio. With this trend, the best fitting ice has an  $\text{H}_2\text{O}:\text{CH}_4$  abundance ratio in the range 2–16. However, the observed width and peak position can also be obtained by adding  $\text{CH}_3\text{OH}$  to mixtures with a larger  $\text{H}_2\text{O}:\text{CH}_4$  abundance ratio (Fig. 2). Thus, the ‘strong interstellar mixture’ ( $\text{H}_2\text{O}:\text{CH}_3\text{OH}:\text{CO}:\text{CH}_4=70:40:1:1$ ) also fits the observations (Fig. 3). Furthermore, the mixture  $\text{CH}_3\text{OH}:\text{CH}_4=30:1$  provides a good fit, when warmed-up to 30 K. We conclude that solid  $\text{CH}_4$  is embedded in a matrix of polar molecules, but no clear distinction between  $\text{H}_2\text{O}$ - and  $\text{CH}_3\text{OH}$ -rich ices and the mixing ratio of these molecules can be made.

A reliable check on the identification of the observed  $7.67 \mu\text{m}$  absorption feature as the deformation mode of interstellar solid  $\text{CH}_4$  can be made by observing the  $3.32 \mu\text{m}$  stretching mode as well. In laboratory experiments, the integrated optical depth of the stretching mode is  $\sim 50\%$  larger than the deformation mode. For polar molecules the width of the stretching mode is also  $\sim 50\%$  larger (Hudgins et al. 1993; B97). Therefore, the peak optical depth of the stretching and deformation modes are expected to be similar. The grating scan of NGC 7538 : IRS9 has an RMS noise on optical depth scale of  $\sim 0.13$  at  $3.32 \mu\text{m}$  (Whittet et al. 1996). This is larger than the expected optical depth (Table 1) of the stretching mode of solid  $\text{CH}_4$ , and thus the assignment of the  $7.67 \mu\text{m}$  feature to interstellar  $\text{CH}_4$  can not be directly confirmed or rejected with the currently available data. Note that observations of the  $\text{CH}_4$  stretching mode require very high signal-to-noise, due to the weakness of protostars at these wavelengths and blending with the broad, deep solid  $\text{H}_2\text{O}$  absorption at  $3.0 \mu\text{m}$ . Nevertheless, it is very important to observe this mode as an independent check on the presence and molecular environment of  $\text{CH}_4$  in ice mantles.

It can readily be seen from Fig. 2 that the broad component toward W 33A can not be explained by a solid state  $\text{CH}_4$  mixture. Ro-vibrational absorption spectra of gaseous  $\text{CH}_4$  show a distinct Q-branch at  $7.66 \mu\text{m}$  and sharp, but weaker lines in the P- and R-branches, for the physical conditions toward W 33A



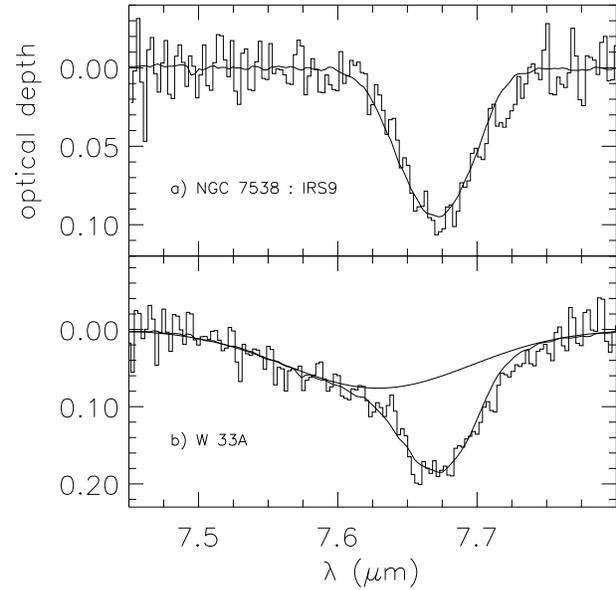
**Fig. 2.** Laboratory spectroscopy of the deformation mode of interstellar  $\text{CH}_4$  compared to the observed peak position and FWHM (cross with error bars). Bullets indicate  $\text{CH}_4$  in a non-polar matrix, circles a polar matrix. The mixing ratio of  $\text{N}_2$ ,  $\text{O}_2$ , and  $\text{CO}_2$  with  $\text{CH}_4$  is 20:1. The ‘S’ indicates the ‘strong interstellar mixture’, and ‘W’ its weak equivalent (Sect. 4; B97). For the other mixtures, a line is drawn between points with an increasing  $\text{CH}_4$  abundance. For each series the highest and lowest abundances relative to  $\text{CH}_4$  are indicated. The temperature for all mixtures is 10 K. Further details can be found in B97.

and at the resolution of our observations (Helmich 1996; B97). The spectrum of W 33A does not show these narrow lines, and therefore we conclude that the broad  $7.63 \mu\text{m}$  absorption feature is not due to gaseous  $\text{CH}_4$ . Instead, this feature may be due to solid  $\text{SO}_2$  (B97).

## 5. Solid $\text{CH}_4$ column densities and chemistry

The column density of ice species can be derived by dividing the integrated optical depth by the molecular band strength  $A$ . We adopt a matrix independent  $A = 7.3 \cdot 10^{-18} \text{ cm molecule}^{-1}$  for the deformation mode of solid  $\text{CH}_4$  (B97). Using this band strength and the Gaussian depths and widths (Table 1), we derive column densities of  $N(\text{CH}_4) = 1.3$  and  $1.7 \cdot 10^{17} \text{ cm}^{-2}$  for NGC 7538 : IRS9 and W 33A respectively.

The column density of solid  $\text{CH}_4$  is compared to other ices in Table 1. There are two independent solid CO components toward these sources: CO in a polar and in a non-polar environment (Tielens et al. 1991). The solid  $\text{CH}_4$  column density is  $\sim 50\%$  of solid CO in a polar environment for both sources. In contrast, while toward the warm source W 33A most of the non-polar CO (sublimation temperature  $\sim 20$  K) has evaporated from the grains, the cold source NGC 7538 : IRS9 has an abundant non-polar CO component. Apparently, there is very little solid  $\text{CH}_4$  associated with this non-polar environment toward NGC 7538 : IRS9. This is consistent with the best fitting laboratory mixtures (Fig. 2). Note that non-polar molecules can be stored in an  $\text{H}_2\text{O}$  matrix up to the  $\text{H}_2\text{O}$  sublimation temperature.



**Fig. 3.** Comparison of the ISO-SWS spectra with the deformation mode of solid  $\text{CH}_4$  in the ‘strong interstellar mixture’ (Sect. 4). **a:** NGC 7538 : IRS9. **b:** W 33A, the laboratory spectrum is added to a Gaussian fit of the underlying broad component, which can not be ascribed to absorption by solid  $\text{CH}_4$ .

There are various ways to form  $\text{CH}_4$  both in the gas phase and on grain surfaces. Each of these processes has its own characteristics which may be ‘traced’ back in the observations. Our observations, as well as the observations by Lacy et al. (1991) indicate a low gas phase  $\text{CH}_4$  column density ( $N(\text{gas})/N(\text{solid}) < 1$ ). This may indicate that the formation of  $\text{CH}_4$  molecules is closely related to interstellar grains.

We can distinguish two classes of models for the formation of interstellar  $\text{CH}_4$  on grains. First,  $\text{CH}_4$  can be formed by hydrogenation of atomic C on grain surfaces. Since solid  $\text{H}_2\text{O}$  is formed in a similar way from accreted oxygen, the predominance of solid  $\text{CH}_4$  in a polar environment is a natural consequence of this model (Tielens & Hagen 1982). The observed low solid  $\text{CH}_4/\text{H}_2\text{O}$  ratio might locate the formation of  $\text{CH}_4$  then deep inside molecular clouds, where most of the elemental carbon is locked up in CO and gaseous C is only a trace species (Keene 1990).

Second, solid  $\text{CH}_4$  can be formed efficiently by UV photolysis of ice mixtures containing  $\text{CH}_3\text{OH}$ , resulting in solid  $\text{CH}_4/\text{CH}_3\text{OH}$  ratios of  $\sim 0.1$  (Allamandola et al. 1988). The observed solid  $\text{CH}_4/\text{CH}_3\text{OH}$  ratio is in good agreement with the laboratory experiments (Table 1; Gerakines et al. 1996). In this model, solid  $\text{CH}_4$  would also be preferentially confined to polar environments. However, the laboratory spectra of photolysed  $\text{CH}_3\text{OH}$  ices also display a narrow feature near  $5.82 \mu\text{m}$ , attributed to formaldehyde ( $\text{H}_2\text{CO}$ ) and higher aldehydes. For irradiation doses consistent with the observed  $\text{CH}_4/\text{CH}_3\text{OH}$  ratio, this band should be  $\sim 50\%$  deeper than the observed  $\text{CH}_4$  feature, but it is not apparent in the spectrum of NGC 7538 : IRS9

**Table 1.** Gaussian parameters of the observed absorption features. The solid CH<sub>4</sub> column density is given as a percentage of solid H<sub>2</sub>O, solid CO in a polar environment (CO-p), solid CO in a non-polar environment (CO-np) and CH<sub>3</sub>OH ice.

Object	$\lambda$ $\mu\text{m}$	FWHM $\mu\text{m}$	$\tau$	N(CH <sub>4</sub> )				
				$10^{17} \text{ cm}^{-2}$	%H <sub>2</sub> O	%CO – p	%CO – np	%CH <sub>3</sub> OH
NGC 7538 : IRS9	7.674±0.003	0.063±0.005	0.092±0.005	1.3±0.1	1.9 <sup>a</sup>	41 <sup>d</sup>	20 <sup>d</sup>	14 <sup>c</sup>
W 33A	7.676±0.005	0.061±0.007	0.120±0.005	1.7±0.2	0.4 <sup>b</sup>	61 <sup>d</sup>	155 <sup>d</sup>	4.4 <sup>c</sup>
	7.631±0.007	0.161±0.010	0.082±0.005	...	...	...	...	...

<sup>a</sup> 3.0  $\mu\text{m}$  band <sup>b</sup> 6.0  $\mu\text{m}$  band <sup>c</sup> 3.54  $\mu\text{m}$  band (Allamandola et al. 1992; Schutte et al. 1996)

<sup>d</sup> Tielens et al. (1991)

(Schutte et al. 1996). Furthermore, the solid CH<sub>4</sub>/CH<sub>3</sub>OH ratio is lower toward W 33A compared to NGC 7538 : IRS9. This anti-correlates with the 4.62  $\mu\text{m}$  absorption band, generally ascribed to energetic processing of ice mantles (Tegler et al. 1995, and references therein). Additional laboratory work is necessary to address these problems and to assess the importance of UV photolysis for the formation of interstellar solid CH<sub>4</sub>.

## 6. Conclusions

We have obtained 7.4–8.0  $\mu\text{m}$  spectra of the highly obscured protostars W 33A and NGC 7538 : IRS9, using the SWS spectrometer on board of the ISO satellite. Both sources exhibit a narrow absorption feature near 7.67  $\mu\text{m}$ . We compare these spectra to laboratory simulations of interstellar solid CH<sub>4</sub> and calculations of gas phase CH<sub>4</sub>. We conclude that the observed 7.67  $\mu\text{m}$  absorption feature toward both sources can be ascribed to absorption by the deformation mode of solid CH<sub>4</sub> in a polar matrix. Both H<sub>2</sub>O-, and CH<sub>3</sub>OH-rich ices can explain the observations. The solid CH<sub>4</sub> column density is 0.4–1.9% of solid H<sub>2</sub>O. No gas phase CH<sub>4</sub> was observed, and therefore we conclude that the formation of CH<sub>4</sub> is closely related to interstellar dust grains. Thus, interstellar CH<sub>4</sub> is likely formed on grain surfaces by hydrogenation of atomic C or by UV photolysis of CH<sub>3</sub>OH in ice mantles. Both models can explain the observed CH<sub>4</sub> abundance and the polarity of the ice in which CH<sub>4</sub> is embedded.

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