

BVRI observations of PSR B0656+14 with the 6-meter telescope

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Abstract. We observed the middle-aged radio pulsar B0656+14 with a CCD detector at the 6-m telescope. Broadband BVRI images show the following magnitudes of the pulsar counterpart: $B = 24.85(+0.19, -0.16)$, $V = 24.90(+0.16, -0.14)$, $R = 24.52(+0.12, -0.11)$, $I = 23.81(+0.27, -0.21)$. We fitted the UV-optical (space + ground-based) data with a two-component model which combines a power law (non-thermal component) with a thermal spectrum emitted by the neutron star surface. The power law component, with the energy power-law index $\alpha = 1.5(+1.1, -1.2)$, dominates in the observed range. Constraints on the thermal component correspond to the Rayleigh-Jeans parameter $G \equiv T_6(R_{10}/d_{500})^2 = 4.1(+2.1, -4.1)$, where $T = 10^6 T_6$ K is the brightness temperature, $R_\infty = 10R_{10}$ km is the neutron star radius as seen by a distant observer, and $d = 500d_{500}$ pc is the distance. The shape of the optical-UV spectrum of PSR B0656+14 differs considerably from those observed from other pulsars: the middle-aged Geminga and young Crab, Vela.

Key words: pulsars: individual: PSR B0656+14 – stars: pre-main sequence – stars: neutron

1. Introduction

Optical radiation has been detected from only $\sim 1\%$ of ~ 750 radio pulsars (see Bignami & Caraveo 1996 for a recent review). Optical counterparts were reported for pulsars of quite different ages, from the very young PSR B0531+21 (Crab pulsar — see Percival et al. 1993) to old PSR B1929+10 and B0950+08 (Pavlov, Stringfellow & Cordova 1996a). Of special interest among these objects are the middle-aged ($\tau \sim 10^5 - 10^6$ yr), soft X-ray pulsars PSR B0656+14, J0633+1746 (Geminga) and B1055–52 studied extensively with the X-ray observatories *ROSAT* and *ASCA*. Their thermal-like X-ray radiation is believed to be emitted by surface layers of the neutron stars (e. g., Ögelman 1995), so that studying this radiation enables one to estimate surface temperatures of neutron stars (NSs) of different ages, to compare them with the current models for

the NS cooling, and to constrain properties of the superdense matter in the NS interiors (e. g., Pavlov et al. 1995). Since the surface temperature, as well as other parameters of radiating layers, can be measured much more accurately if the surface radiation is also observed in the optical–UV range (Pavlov et al. 1996b), investigation of the optical–UV radiation of these objects is potentially very important.

In this paper we report multi-color CCD observations of the middle-aged PSR B0656+14 ($\tau = P/2\dot{P} = 1.1 \times 10^5$ yr) with the 6-meter telescope (BTA) at the Special Astrophysical Observatory of the Russian Academy of Sciences (SAO RAS). The optical counterpart of this pulsar was first detected with the ESO NTT and 3.6-meter telescopes in the V band ($V \sim 25$) by Caraveo et al. (1994a) at a 3σ level. Pavlov et al. (1996a) observed the pulsar with the Faint Object Camera (FOC) onboard the Hubble Space Telescope (*HST*) in a longpass filter F130LP ($\lambda\lambda = 2310 - 4530$ Å). They confirmed the optical detection at a very high confidence level ($S/N = 52$) and concluded that the observed radiation is mainly of a nonthermal (most likely, magnetospheric) origin, although a contribution of a thermal (Rayleigh–Jeans) component emitted by the NS surface cannot be excluded. In order to separate the thermal component, and to understand the nature of the nonthermal radiation, the flux of the object should be measured in several spectral bands covering as broad spectral range as possible. Mignani, Caraveo & Bignami (1997) observed PSR B0656+14 with the *HST* Planetary Camera (PC) in the F555W filter and determined its V magnitude more accurately: $V = 25.1 \pm 0.1$. Obviously, observations with other filters are needed to find the shape of the spectral flux of the optical counterpart. Thus, we included such observations in our general program of deep broad-band CCD photometry of nearby pulsars of the northern sky with the 6-meter telescope.

In our first BTA observations of PSR B0656+14 (January 1996) in the B , V and R filters (Kurt et al. 1997a) we obtained estimates of its brightness in the V and, for the first time, B and R bands and confirmed the nonthermal nature of its optical spectrum. However, because of poor weather conditions during those observations, the fluxes were determined with substantial uncertainties. Subsequent observations of the same object (November 1996) were more successful (Kurt et al. 1997b).

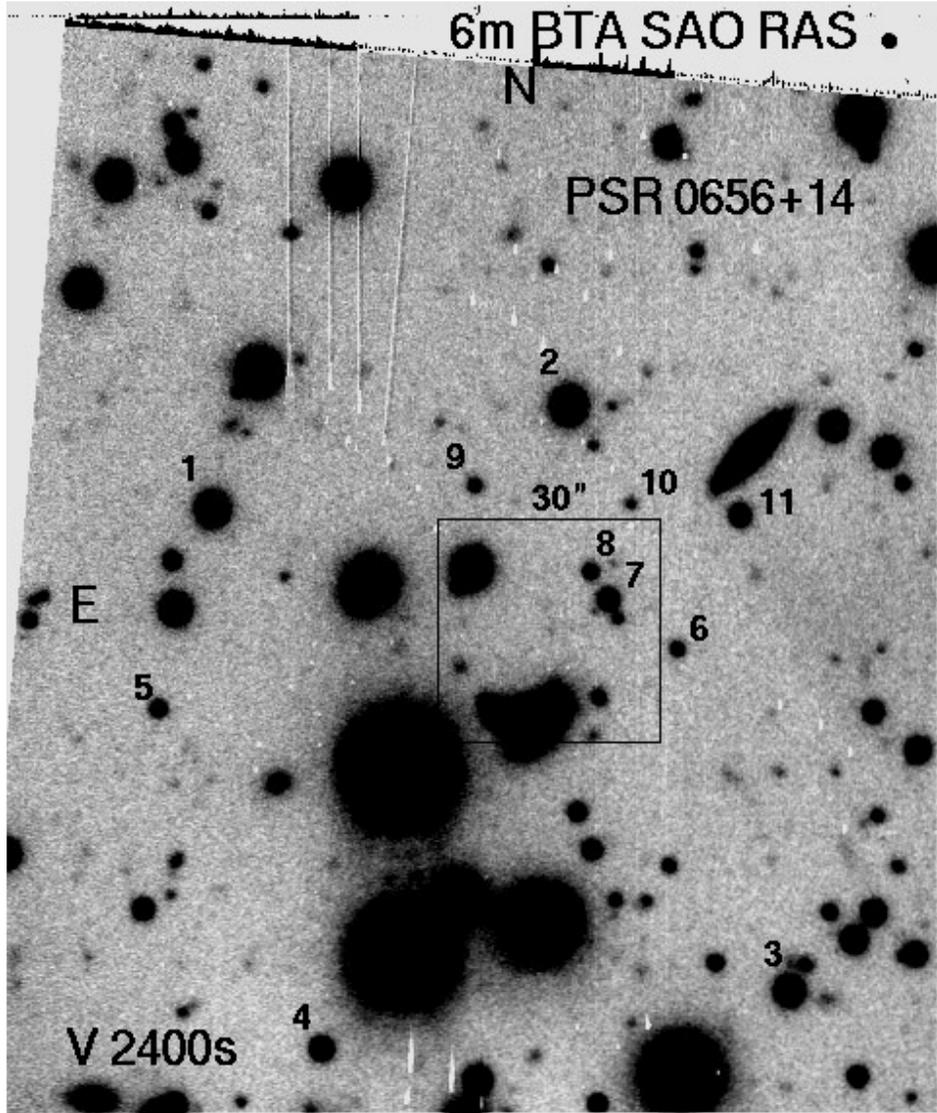


Fig. 1. Image of the PSR B0656+14 field in the V filter with accumulation time of 2400 sec. The center of the $30'' \times 30''$ fragment coincides with the pulsar position. The numbers denote the stars whose magnitudes are given in Table 3.

We describe these observations and their reduction in Sect. 2 and discuss the properties and origin of the optical spectrum of PSR B0656+14 in Sect. 3.

2. Observations and data analysis

2.1. Observations

Photometric observations of PSR B0656+14 were carried out with the 6-meter telescope at SAO RAS on 11/12 November, 1996 with a CCD photometer installed at the Primary Focus. We used the “Electron ISD017A” CCD; its format of 1040×1160 pixels corresponds to the field of view $2'.38 \times 2'.66$. The CCD photometer was employed in the 2×2 binning mode, so that each of the 520×580 zoomed pixels (referred to as ‘pixels’ hereafter) has the angular size of $0''.274 \times 0''.274$. We used the gain of $2.3e^-$ per DN (Data Number) pixel. The shutter timing accuracy is better than 0.1 sec.

The observations were carried out with filters close to the B , V , R and I filters of the Cousins system. Table 1 gives basic parameters of observational conditions: starting time of each exposure, filter and exposure number, duration of exposure, seeing, zenith distance Z , and sky background in stellar magnitudes per arcsec².

2.2. Data reduction

Standard data reduction includes subtraction of the “bias”, — an additional component of the CCD signal, flat-fielding (correction for non-uniform sensibility of the detector elements), and removing of space particle traces. The data were processed with the use of the MIDAS software. Processing involved all the obtained images except for R*4 where the space particle trace was in the immediate vicinity of the pulsar position. Fig. 1

Table 1. Observations of PSR B0656+14 on November 11/12, 1996.

time UT	exp. ID	exposure (sec)	seeing (arcsec)	Z (deg)	backgr. (mag/ arcsec ²)
22:37	B 1	600	1.64	40.63	22.41
22:49	V 1	600	1.56	39.00	21.60
23:01	R 1	600	1.42	37.43	20.70
23:11	I 1	300	1.31	35.95	19.39
23:18	I 2	300	1.36	35.17	19.34
23:25	B 2	600	1.69	34.42	22.39
23:36	V 2	600	1.42	33.21	21.61
23:47	R 2	600	1.31	32.13	20.80
23:59	I 3	300	1.42	31.22	19.52
00:06	I 4	300	1.33	30.76	19.47
00:12	I 5	300	1.26	30.39	19.45
00:19	B 3	600	1.67	30.06	22.34
00:30	V 3	600	1.31	29.65	21.52
00:42	R 3	600	1.45	29.43	20.64
01:13	R*4	600	1.45	29.88	20.64
01:25	I 6	300	1.37	30.42	19.29
01:31	I 7	300	1.40	30.50	19.32
01:41	I 8	300	1.30	31.57	19.41
01:48	B 4	600	1.75	32.10	22.27
01:59	V 4	600	1.61	33.18	21.41
02:11	R 5	600	1.51	34.39	20.57
02:22	R 6	600	1.53	35.74	20.53
02:34	I 9	300	1.51	37.20	19.22
02:41	I 10	300	1.53	38.05	19.05
02:47	I 11	300	1.51	38.95	18.75

Table 2. Coordinates of objects in the PSR B0656+14 field.

No.	α_{1950}	δ_{1950}	relative coordinates (arcsec)	distance from PSR (arcsec)
1	06:57:00.97	14:18:51.5		
2	06:56:57.62	14:19:05.1		
3	06:56:55.63	14:17:43.6		
4	06:57:00.02	14:17:36.6		
5	06:57:01.51	14:18:24.1		
6	06:56:56.63	14:18:31.3		
o1,PSR	06:56:57.84	14:18:34.2	206.5; 203.0	0.0
o2	06:56:57.75	14:18:36.2	207.5; 204.3	1.7
o3	06:56:57.87	14:18:37.4	206.1; 207.1	4.1
o4	06:56:57.47	14:18:34.7	212.0; 204.0	5.1
o5	06:56:57.58	14:18:34.7	209.5; 204.0	3.1

shows the field around the pulsar (sum of 4 exposures in the V filter).

2.3. Astrometrical referencing

According to Thompson and Cordova (1994), the VLA position of the pulsar is $\alpha_{1950} = 06^h 56^m 57^s.942$; $\delta_{1950} = 14^\circ 18' 33''.80$ (epoch 1992.98); the formal uncertainties of the radio position ($0''.09$ and $0''.03$ for α and δ , respectively) do not include possible systematic errors which may be as large as $\sim 0''.2$ (cf. Pavlov et al. 1996a). With allowance for the proper motion of the pulsar,

$\mu_\alpha = +0.73 \pm 20$ and $\mu_\delta = -26 \pm 13$ milliarcsec yr⁻¹ (Pavlov et al. 1996a), the expected position of the pulsar at the epoch of our observations (1996.86) is $\alpha_{1950} = 06^h 56^m 57^s.96(\pm 0''.02)$, $\delta_{1950} = 14^\circ 18' 33''.7(\pm 0''.2)$.

The position of the candidate optical counterpart relative to nearest bright stars found with the use of the *HST* Guige Star Catalog (GSC) has been presented by Caraveo et. al. (1994b). The astrometric referencing of our data is based on coordinates taken for the field stars around the pulsar from the Digitized Sky Survey (nominal accuracy $\sim 1''$).

Table 2 gives the coordinates of the pulsar optical candidate, 6 bright stars (denoted by the numbers 1 through 6 in Fig. 1), and 4 objects within $\sim 5''$ around the pulsar. For these objects, the coordinates corresponding to the frames of Figs. 3, 4 and 6 are also presented.

2.4. Photometry

Photometric referencing was carried out by observing four standard stars from the Landolt catalog (1992): 112822, GD71 and RU 149 (8 stars) and star 2015 from Weis (1991). Photometric conditions remained stable during the whole night, with an insignificant increase of absorption at the blue part of the spectrum at the end of the observations. For the extinction factors, we used average values of the observatory: $k_b = 0.34$, $k_v = 0.21$, $k_r = 0.15$, $k_i = 0.10$ (in stellar magnitudes). From observations of the standard stars we derived the following transition equations from the instrument system (magnitudes b, v, r, i) to the Cousins system (magnitudes B, V, R, I):

$$\begin{aligned}
 B - b &= 26.14 + 0.16(b - v) \\
 V - v &= 26.30 - 0.11(b - v) \\
 R - r &= 26.62 - 0.02(v - r) \\
 I - i &= 25.77 + 0.07(r - i) .
 \end{aligned}
 \tag{1}$$

The instrumental magnitudes were calculated as

$$b = -2.5 \log(F_b/t_{\text{exp}}) - \Delta_a m_b - k_b / \cos Z , \tag{2}$$

etc, where F is the DN value in a given aperture, t_{exp} is the exposure time, $\Delta_a m$ is the correction for the finite aperture (obtained from the Point Spread Function [PSF] of bright stars), and $k / \cos Z$ is the correction for the atmospheric extinction. Errors of the numerical coefficients in Eq. (1) do not exceed ± 0.01 .

The errors of the stellar magnitudes can be determined by measuring the signal-to-noise (S/N) ratios:

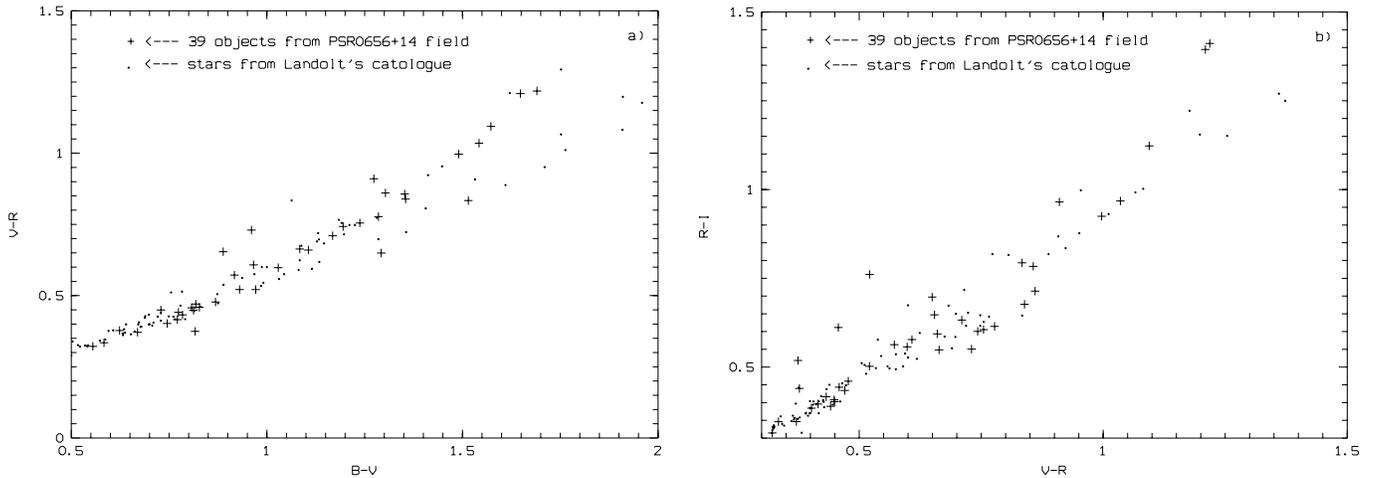
$$\Delta m = -2.5 \log\left(1 \pm \frac{1}{S/N}\right), \tag{3}$$

$$\frac{S}{N} = \frac{F\sqrt{Gain}}{\sqrt{F + Sky \times N_{\text{pix}}}} , \tag{4}$$

where N_{pix} is the number of pixels in the aperture chosen, Sky is the value of the local background in DN per pixel, and $Gain =$

Table 3. Stellar magnitudes of 11 stars around the pulsar.

	<i>B</i> ($\lambda_{\text{eff}} = 4448 \text{ \AA}$, FWHM=1008 \AA)	<i>V</i> ($\lambda_{\text{eff}} = 5505 \text{ \AA}$, FWHM=827 \AA)	<i>R</i> ($\lambda_{\text{eff}} = 6588 \text{ \AA}$, FWHM=1568 \AA)	<i>I</i> ($\lambda_{\text{eff}} = 8060 \text{ \AA}$, FWHM=1542 \AA)
1	19.35 \pm 0.01	18.31 \pm 0.01	17.72 \pm 0.01	17.17 \pm 0.01
2	18.71 \pm 0.01	18.02 \pm 0.01	17.66 \pm 0.01	17.32 \pm 0.01
3	19.90 \pm 0.01	18.69 \pm 0.01	17.96 \pm 0.01	17.36 \pm 0.01
4	20.94 \pm 0.02	19.84 \pm 0.01	19.19 \pm 0.01	18.64 \pm 0.01
5	22.61 \pm 0.02	21.31 \pm 0.01	20.54 \pm 0.01	19.93 \pm 0.01
6	22.70 \pm 0.03	21.96 \pm 0.02	21.52 \pm 0.01	21.12 \pm 0.02
7	21.99 \pm 0.03	20.31 \pm 0.02	19.13 \pm 0.01	17.74 \pm 0.02
8	23.08 \pm 0.03	21.52 \pm 0.02	20.50 \pm 0.01	19.53 \pm 0.02
9	23.58 \pm 0.03	22.06 \pm 0.02	21.24 \pm 0.01	20.44 \pm 0.02
10	24.52 \pm 0.03	22.88 \pm 0.02	21.92 \pm 0.01	21.08 \pm 0.02
11	21.38 \pm 0.03	20.40 \pm 0.02	19.89 \pm 0.01	19.39 \pm 0.02

**Fig. 2.** **a** $V - R$ vs. $B - V$ diagram for stars of the PSR B0656+14 field. Crosses denote the color indices for the field objects. Points show the color indices for stars from the Landolt catalogue. **b** $R - I$ vs. $V - R$ diagram for the field stars.

$2.3e^-/\text{DN}$ is the gain. The CCD readout noise is $\sim 8e^-$ or 4–5 DN and one was negligible in our observations.

Making use of Eqs. (1)–(4), we measured the magnitudes of relatively bright stars of the pulsar field (individual images were summed for each filter). The Cousins magnitudes of 11 of these stars (denoted by the numbers 1 through 11 in Fig. 1), together with the effective wavelengths and widths of the Cousins filters, are presented in Table 3. Fig. 2 shows the comparison of the color indices for 39 objects of the pulsar field, from $V = 17.5$ to $V = 22.5$, with the color indices for stars from the Landolt catalog. The color-color dependences are close to each other, which attests both low absorption in the Galaxy for the direction towards PSR B0656+14 and the absence of systematic shifts in our photometric system.

2.5. Optical counterpart of PSR B0656+14

Contour plots of the pulsar vicinity in all the 4 filters obtained from smoothed (3×3 pixels) images are shown in Figs. 3 and 4. The first five isophotes (for lower brightnesses) correspond to the following values above the background:

$Sky + n\sqrt{Sky/N_{\text{pix}}}$, with $n = 1, 2, 3, 4, 5$, where $N_{\text{pix}} = 28$ corresponds to the aperture diameter of 6 pixels ($1''.6$), and the value of Sky is that for the background near the pulsar. The isophotes for higher brightnesses are scaled arbitrarily to demonstrate brighter details. It is seen from Figs. 3 and 4 that 5 objects are detected within $\lesssim 5''$ around the expected pulsar position. The position of one of them (“o1,PSR” in Table 2 and Fig. 6) coincides, within the budget of astrometry errors, with the radio position, so that this object is the most plausible candidate for the optical counterpart of PSR B0656+14. The nearest object “o2” (detected in the I filter only) lies to the North at a distance of $\approx 2''$. Since the objects in the near vicinity are not seen simultaneously in all 4 filters, we have shown them, for the sake of convenience, in the separate scheme of Fig. 6.

The photometry of the pulsar candidate (Table 4) was carried out for several aperture radii R_a , to optimize the S/N ratio. The object “o2”, whose brightness in the I filter is comparable to that of the pulsar candidate, starts to influence the photometry of the candidate at $R_a \gtrsim 4$ pixels ($\gtrsim 1''$). Most notations of Table 4 have been described above. The columns m_{ins} and m_{cous} are for the magnitudes in the instrumental and Cousins systems,

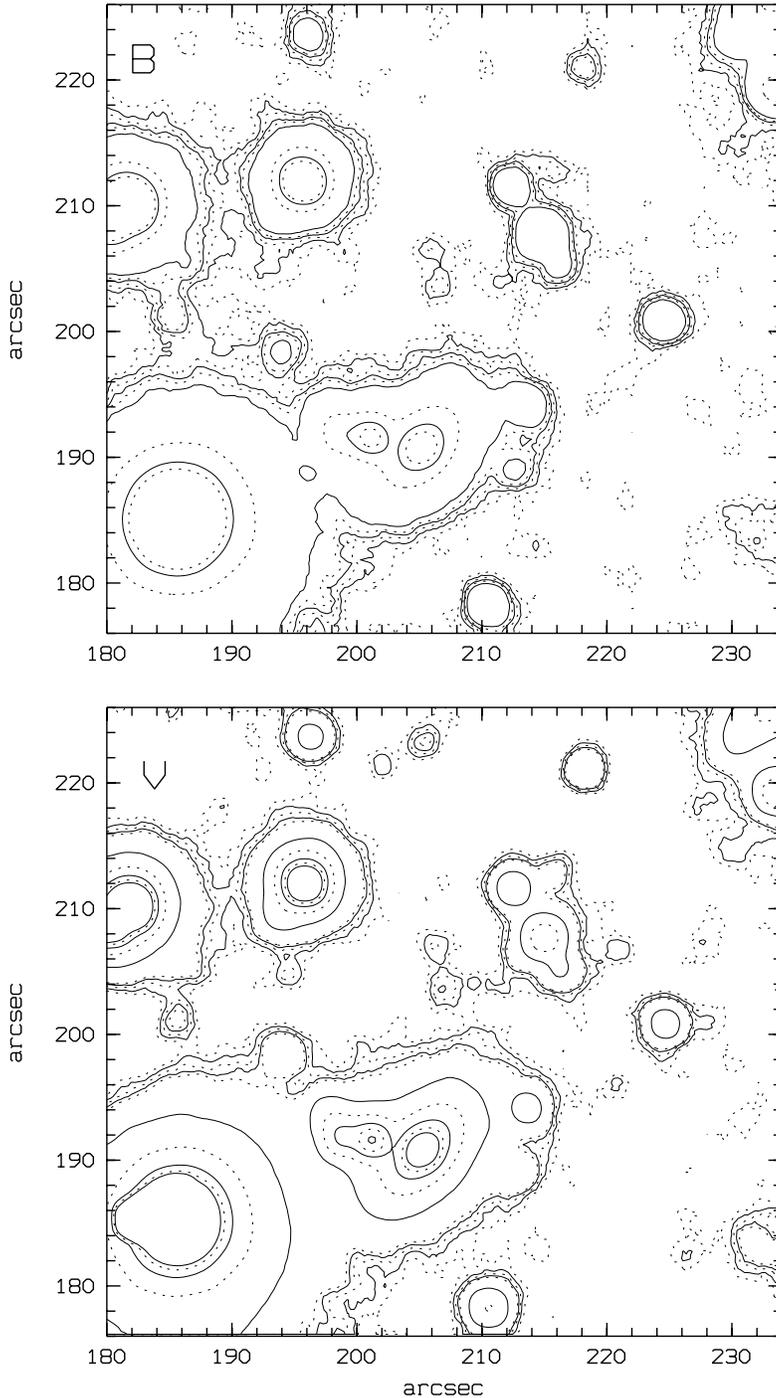


Fig. 3. Contour plots of the PSR B0656+14 vicinity in the *B* and *V* filters. The size of each fragment is $50'' \times 54''$.

respectively. The last column shows the source flux (in 10^{-29} erg cm^{-2} s^{-1} Hz^{-1}) obtained from the Cousins magnitudes with the aid of the data on α Lyr by Fukigita et al. (1995)

2.6. Objects in near vicinity of the PSR B0656+14 counterpart

The coordinates of the 5 objects within $\lesssim 5''$ of the expected pulsar position (see Figs. 5 and 6) are presented in Table 2.

The object “o1,PSR”, whose brightness distribution is consistent with that of a point source, is the only probable candidate

for the pulsar counterpart, based upon both its position and photometry.

The object “o3” is seen in all our *B, V, R, I* images, as well as in the contour plots of Fig. 2 of Caraveo et al. (1994a). It has been also detected as a faint extended object in the *HST* F130LP image (Pavlov et al. 1996a). Moreover, it can be seen even in our first *B* images (Kurt et al. 1997a) as a background enhancement exactly at the same position. Thus, the presence of this faint (extended?) object at $\simeq 4''$ to the North of the pulsar is firmly established. The object “o4” has also been detected with

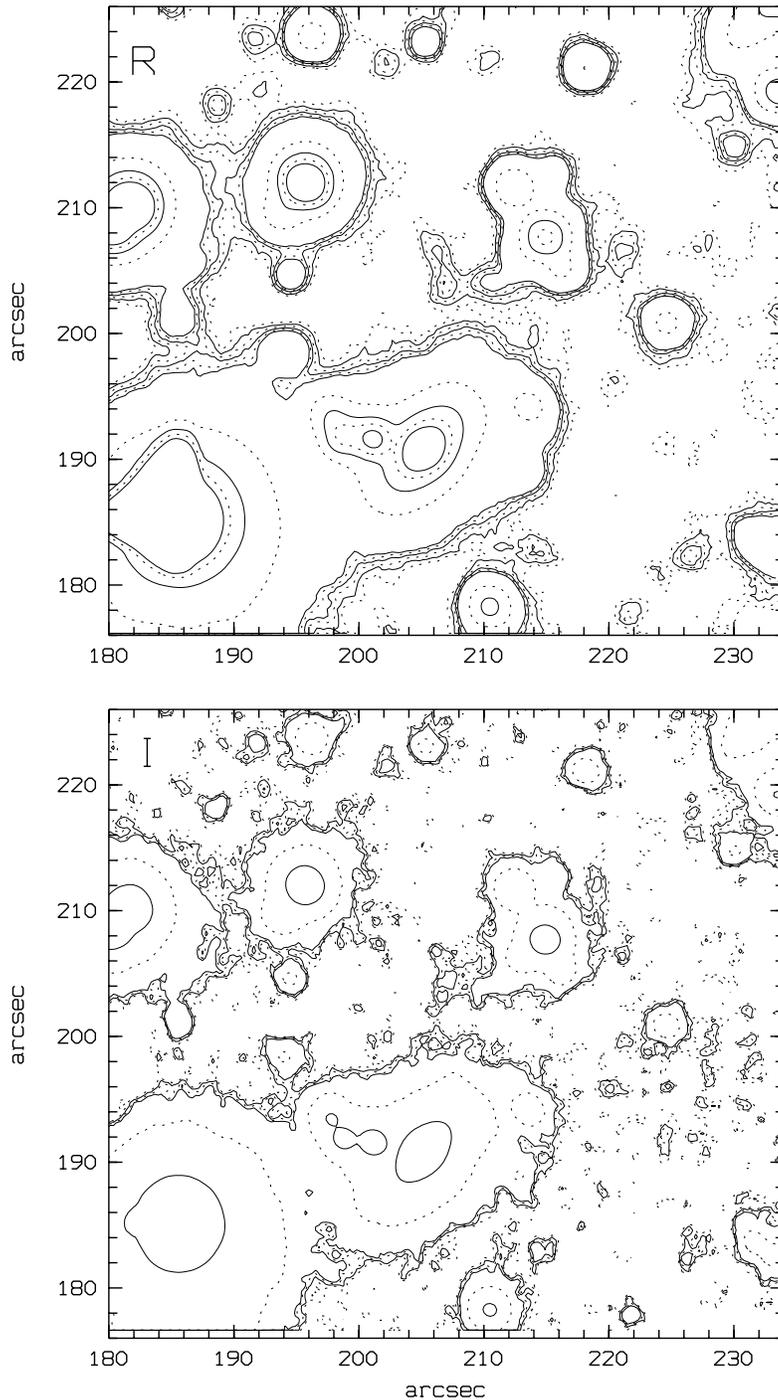


Fig. 4. Contour plots of the PSR B0656+14 vicinity in *R* and *I* filters. The size of each fragment is $50'' \times 54''$.

the ESO-3.6m and NTT telescopes in the *V* filter (Caraveo et al. 1994a), although it is difficult to distinguish between the objects “o4” and “o5” because of their obviously extended nature. These objects are outside of the $7''.4 \times 7''.4$ frame of the *HST* FOC image presented by Pavlov et al. (1996a).

The object “o2”, closest to the pulsar, is reliably detected only in the *I* filter. In other filters only a non-significant enhancement of the background is seen at this position. We suppose that the object “o2” is not seen in the *HST* FOC image because it is too red to be detected in the F130LP band.

Table 5 gives estimates for magnitudes of these objects. These estimates and the distances of the objects relative to “o1,PSR” (Table 2) allow us to conclude that these objects do not affect the measured values of the *B*, *V*, *R* magnitudes of the pulsar counterpart. A minor effect about 3% from whole flux of the nearest object “o2” on the *I* magnitude determined for the pulsar counterpart with the aperture radius of 3 pixels (diameter of $1''.6$) can be expected, but it should not exceed the errors of our measurement.

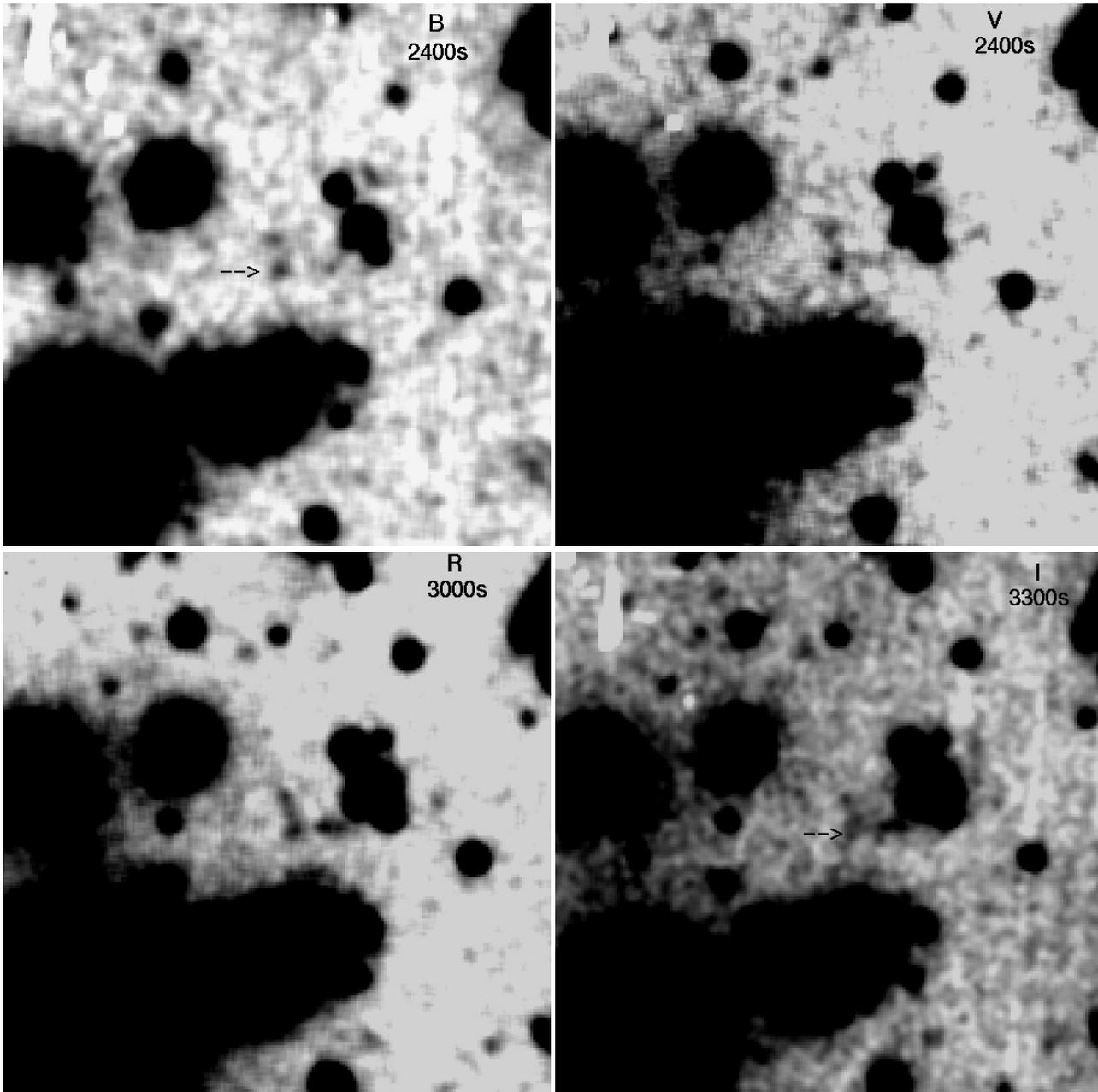


Fig. 5. Vicinity of the pulsar PSR B0656+14 in the *B*, *V*, *R*, *I* filters. The size of each fragment is $54'' \times 54''$. The optical counterpart of the pulsar (“o1,PSR”, cf. Fig. 6) is indicated by arrows in the *B* and *I* images. The nearest object (“o2” in Fig. 6) lies to the North at a distance of $\approx 2''$.

3. Discussion.

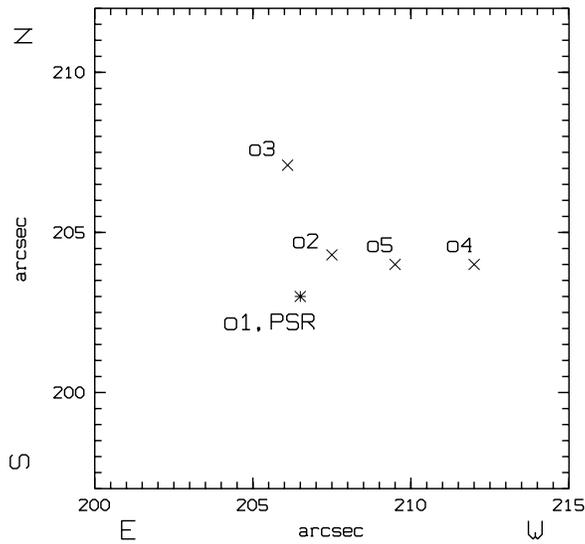
Fig. 7 shows the UV-optical energy fluxes in different bands for the PSR B0656+14 counterpart obtained with the *HST* and 6-m telescope. It demonstrates that our results (crosses *B*, *V*, *R* and *I*), being compatible with previously obtained data, considerably extend redward the observed frequency domain, and provide more accurate fluxes in the *V* and *B* bands. The overall spectrum firmly confirms a nonthermal origin of the optical radiation suggested by Pavlov et al. (1996a). On the other hand, the relatively high value of the F130LP flux makes very plau-

sible that the thermal radiation from the NS surface contributes noticeably into the UV part of the spectrum. Parameters of the thermal and nonthermal fluxes can be constrained by fitting the observed spectrum with a two-component model (cf. Pavlov et al. 1996a), a sum of the power-law and thermal spectra (the latter has the form of a Rayleigh-Jeans spectrum at these frequencies for plausible NS temperatures):

$$f(\nu) = \left[f_0 \left(\frac{\nu}{\nu_0} \right)^{-\alpha} + g_0 \left(\frac{\nu}{\nu_0} \right)^2 \right] \times 10^{-0.4A(\nu)}, \quad (5)$$

Table 4. Photometry of the PSR B0656+14 optical counterpart.

filter	Sky/t_{exp} (DN/pix/s)	R_a (pix)	N_{pix}	S/N	$-2.5 \log(F/t_{exp})$	$\Delta_a m$	$k/\cos Z$	m_{ins}	m_{cous}	Flux (μJy)
<i>B</i>	2.465	3	28	5.6	+0.13	1.01	0.43	-1.30	24.85(+0.19, -0.16)	0.46 ± 0.08
		4	50	6.1	-0.24	0.63				
		5	79	6.1	-0.42	0.41				
<i>V</i>	6.429	3	28	6.6	-0.29	0.84	0.26	-1.35	24.90(+0.16, -0.14)	0.40 ± 0.06
		4	50	6.7	-0.58	0.51				
		5	79	7.1	-0.73	0.33				
<i>R</i>	18.207	3	28	7.9	-1.11	0.81	0.19	-2.09	24.52(+0.12, -0.11)	0.47 ± 0.06
		4	50	9.8	-1.34	0.48				
		5	79	9.3	-1.60	0.30				
<i>I</i>	29.922	3	28	4.6	-1.02	0.81	0.12	-1.95	23.81(+0.27, -0.21)	0.71 ± 0.16
		4	50	4.2	-1.42	0.49				
		5	79	4.1	-1.67	0.31				

**Fig. 6.** Objects from nearest vicinity of the pulsar are indicated whose coordinates are given in Table 2, and estimates of stellar magnitudes are given in Table 5. Coordinates (in arcseconds) correspond to those in Figs. 3 and 4.**Table 5.** Stellar magnitudes of objects around PSR B0656+14.

object	<i>B</i>	<i>V</i>	<i>R</i>	<i>I</i>
o2	-	-	-	23.83 (+0.21; -0.19)
o3	25.0 ± 0.2	25.4 ± 0.4	24.63 (+0.15; -0.13)	23.93 (+0.24; -0.20)
o4	25.5 ± 0.35	25.5 ± 0.5	24.59 (+0.21; -0.17)	23.50 (+0.12; -0.10)
o5	-	-	24.35 (+0.14; -0.12)	23.54 (+0.20; -0.17)

where ν_0 is an arbitrary reference frequency (we chose $\nu_0 = 8.766 \times 10^{14}$ Hz, which corresponds to $\lambda_0 = 3420$ Å), $A(\nu)$ is the interstellar extinction (see, e. g., Savage & Mathis 1979), f_0 and g_0 are the values of the nonthermal and thermal fluxes

at $\nu = \nu_0$. The latter can be expressed, for the chosen value of ν_0 , as

$$g_0 = 3.116 \times 10^{-31} G \frac{\text{erg}}{\text{cm}^2 \text{ s Hz}}, \quad G = T_6 \left(\frac{R_{10}}{d_{500}} \right)^2, \quad (6)$$

where $T = 10^6 T_6$ K is the NS brightness temperature, $R_\infty = 10 R_{10}$ km is the NS radius (as seen by a distant observer), and $d = 500 d_{500}$ pc is the distance to the pulsar. According to Pavlov et al. (1996a), a plausible value for the color excess, which determines the extinction curve $A(\nu)$, is $E(B - V) = 0.03$ ($A_V = 0.09$). We fixed $E(B - V)$ at this level and fitted the fluxes in the 6 bands (*B*, *V*, *R*, *I*, F555W and F130LP) with Eq. (5), varying f_0 , α and G . The best-fit model (minimum $\chi^2 = 4.69$ for 3 degrees of freedom) corresponds to $f_0 = 1.75 \times 10^{-30}$ erg cm $^{-2}$ s $^{-1}$ Hz $^{-1}$, $\alpha = 1.53$ and $G = 4.1$. The best-fit spectrum is shown in Fig. 7 with the solid line, and the best-fit thermal and power-law components with dashed lines.

The value of G can be compared with what is expected from extrapolations of the soft X-ray spectral flux observed with the *ROSAT* and *ASCA* (Finley, Ögelman & Kiziloğlu 1992; Anderson et al. 1993; Greiveldinger et al. 1997) to the optical range. The result of this extrapolation depends on the model assumed to fit the X-ray spectrum. The blackbody fit of the *ROSAT* spectrum yields $T_6 = 0.9 \pm 0.04$ for $R_{10}/d_{500} = 1.0 \pm 0.2$ (Finley et al. 1992), which corresponds to the Rayleigh-Jeans parameter $G = 0.9 \pm 0.3$. A fit of the same data with a set of hydrogen NS atmosphere models (Pavlov et al. 1995) gives, for $B = 4.7 \times 10^{12}$ G, the radius-to-distance ratio $R_{10}/d_{500} = 2.3 \pm 0.4$ and the effective temperature $T_\infty^{\text{eff}} = (5.3 \pm 0.5) \times 10^5$ K (Anderson et al. 1993). Since the model spectrum differs from the blackbody spectrum, there is no a universal relation between the effective and brightness temperatures; however, the model spectrum in the optical range closely reminds, and can be fitted with, the Rayleigh-Jeans spectrum with $T_6 \simeq 0.9 T_\infty^{\text{eff}}$, thus giving $G \simeq 2.6 \pm 0.7$, almost thrice greater than G inferred from the blackbody model.

We see that the best-fit value of G obtained from the available optical-UV observations exceeds the predictions of both the blackbody and hydrogen atmosphere interpretations of the soft

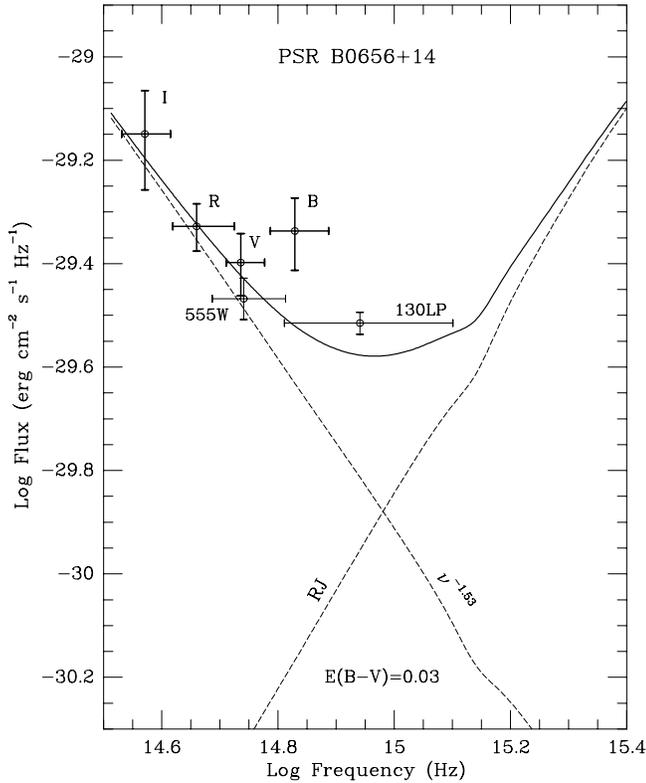


Fig. 7. Fluxes observed from the PSR B0656+14 optical counterpart in different spectral bands. The crosses B , V , R and I show our results; the crosses 130LP and 555W show the fluxes detected with the *HST* FOC/F130LP by Pavlov et al. (1996a) and WFPC2/F555W by Mignani et al. (1997). Dashed crosses show the fluxes obtained by Caraveo et al. (1994; V band) and Kurt et al. (1997a; B band). The 3σ upper limits on the flux from undetected sources in the B, V, R, I filters are shown for the aperture diameter of 6 pixels ($\approx 1''.6$). The solid line shows the best fit of the fluxes in the 6 bands with the two-component model (power law + Rayleigh-Jeans spectrum; Eq. [5]) for the color excess $E(B-V) = 0.03$. The dashed lines are the components of this model.

X-ray flux. However, as Fig. 8 shows, the fitting parameter G is constrained poorly by our data: $0 < G < 6.2$ at a 1σ level. The main reason of that is the lack of data in the far-UV range where the thermal component gives the main contribution. Thus, we can only conclude that the results are consistent with the presence of a thermal component, but we cannot distinguish between different NS surface models and can only estimate upper limits for the Rayleigh-Jeans parameter, $G < 6.2$, 6.8 and 7.3 at the confidence levels of 68, 90 and 99%. We expect that better constraints will be obtained from recent *HST* FOC observations of this pulsar in the UV range (Pavlov, Welty & Córdoba 1997).

Fig. 7 clearly demonstrates that the pulsar radiation in the IR-optical range is predominantly of a nonthermal origin, in agreement with the conclusion of Pavlov et al. (1996a). An independent confirmation of the nonthermal nature comes from the high pulsed fraction observed by Shearer et al. (1997) in the B band. Although the slope of the nonthermal spectrum is not well constrained because of relatively large errors (especially in the I band): $0.3 < \alpha < 2.6$ at a 1σ level, it looks

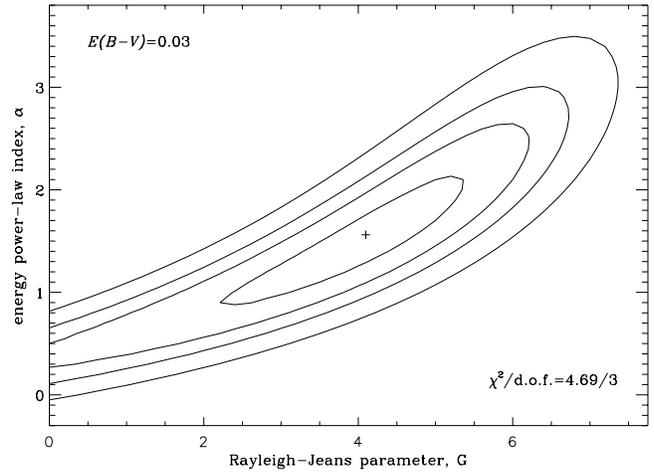


Fig. 8. Confidence levels (20, 68, 90 and 99%) for fitting of the observed spectrum of PSR B0656+14 with the two-component model (Eq. [5])

much steeper than those of the (much younger) Crab pulsar, $\alpha = -0.11 \pm 0.13$ (Percival et al. 1993; see also Ransom et al. 1994). The nonthermal component of the younger Vela pulsar has the slope $\alpha \approx 0.3$ in the VBU range, but its R flux is about twice lower than the extrapolation of the power law (Nasuti et al. 1997). Even more drastic difference is seen between the shapes of the optical spectra of PSR B0656+14 and the older pulsar Geminga ($\tau = 3 \times 10^5$ yr) which looks very much alike B0656+14 in the soft X-ray range. Geminga's flux has a minimum at the B band, is peaked at the V band, and falls off sharply redward, which allowed Bignami et al. (1996) to suggest the presence of a proton cyclotron line centered at $\lambda \approx 6000$ Å. On the contrary, the flux of PSR B0656+14 grows redward of V , thus demonstrating that the properties of the nonthermal IR-optical-UV radiation are neither the same in different pulsars nor they have a clear correlation with the pulsar age.

The apparent excess of the B flux with respect to acceptable continuum fits (Fig. 7), responsible for the relatively large value of the minimum χ^2 , may indicate the presence of a spectral feature at these wavelengths, perhaps similar to that reported for Geminga. More observations are needed to verify whether the feature is real. But at least three broad emission lines above 3600 Å were recently detected indeed in a prism spectrum of PSR B0540-69 with the FOS/*HST* by Hill et al. (1997). The authors connect directly the presence of these strong emissions with a nebula surrounding PSR B0540-69.

The nonthermal component has also been observed in high-energy radiation from PSR B0656+14. Greiveldinger et al. (1997) reported a power-law component in the X-ray spectrum of PSR B0656+14, with $\alpha = 0.5 \pm 1.1$, which dominates at $E \gtrsim 2$ keV. Ramanamurthy et al. (1996) presented an evidence for γ -ray emission ($E > 100$ MeV), with $\alpha = 1.8 \pm 0.3$. Being extrapolated to the X-ray range, our best-fit power-law component ($\alpha = 1.53$) gives the X-ray flux about 3 orders of magnitude lower than observed by Greiveldinger et al., i. e., the nonthermal flux should grow somewhere between the UV and X-ray ranges. To merge with the X-ray power-law component at

$E \sim 2$ keV, our optical flux should have a slope $\alpha \sim 0.3 - 0.4$, hardly compatible with our fit (see Fig. 8), and in this case the spectrum should steepen at higher energies in order to be compatible with the γ -ray data. It is, however, hard to believe that the spectrum would have the same slope in the optical and high-energy ranges — in fact, none of the four other pulsars with investigated nonthermal optical components shows such behavior.

To put more stringent constraints on the parameters of the thermal and nonthermal components, more observations are needed. Observations in the IR range would be most useful for elucidating the slope α of the nonthermal component, whereas far-UV observations would be crucial to measure the Rayleigh-Jeans parameter G proportional to the temperature of the NS surface. Since G and α are correlated with each other (see Fig. 8), measuring one of these parameters with a higher accuracy would immediately constrain the other parameter. It would be also useful to obtain more detailed spectrum around the B band in order to understand whether the excess we observed is associated with a spectral line. Such observations are quite feasible with both the *HST* and large ground-based telescopes.

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