

*Letter to the Editor***ORFEUS II echelle spectra: Absorption by H₂ in the LMC***K.S. de Boer¹, P. Richter¹, D.J. Bomans^{2,**}, A. Heithausen³, and J. Koornneef⁴¹ Sternwarte, Universität Bonn, Auf dem Hügel 71, D-53121 Bonn, Germany² Astronomy Department, University of Illinois at Urbana-Champaign, 1002 West Green St., Urbana, IL 61801, USA³ Radioastronomisches Institut, Universität Bonn, Auf dem Hügel 71, D-53121 Bonn, Germany⁴ Kapteyn Institute, Postbus 800, 9700 AV Groningen, The Netherlands

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Abstract. We report the first detection of H₂ absorption profiles of LMC gas on the line of sight to star 3120 in the association LH 10 near the emission nebula N 11B. Transitions found in the Lyman band are used to derive a total column density $N(\text{H}_2) = 6.6 \times 10^{18} \text{ cm}^{-2}$. Excitation temperatures of $\leq 50 \text{ K}$ for levels $J \leq 1$ and of $\simeq 470 \text{ K}$ for levels $2 \leq J \leq 4$ of H₂ are derived. We conclude that moderate UV pumping influences the population even of the lowest rotational states in this LMC gas.

Key words: space vehicles – ISM: molecules – galaxies: ISM – Magellanic Clouds: LMC – stars: individual: LH 10:3120 – ultraviolet: ISM

1. Introduction

The last instrument capable of high resolution spectroscopy in the 900–1200 Å range was the *Copernicus* satellite (Spitzer et al. 1973), working in the 1970s. This part of the spectrum contains absorption transitions of the Lyman and Werner Bands of molecular hydrogen, H₂, and of transitions of O VI and other species highly relevant for interstellar medium studies.

The *ORFEUS* (Krämer et al. 1990) and *IMAPS* (Jenkins et al. 1988) experiments on the *ASTRO-SPAS* space shuttle platform has provided access to the far UV spectral range in great detail. The *ORFEUS* telescope feeds two spectrographs. The Heidelberg-Tübingen echelle gives spectra from 912 to 1410 Å with $\lambda/\Delta\lambda \leq 10^4$ which are recorded with a microchannel plate detector. Good S/N is achieved in exposure times of the order of 1 hr of hot objects with $V < 12 \text{ mag}$ (Krämer et al. 1990). The UCB spectrograph produces spectra over the range of 390 to 1220 Å (Hurwitz et al. 1998) with a resolution too low for detailed interstellar work. *IMAPS* has its own telescope, works at $\lambda/\Delta\lambda \sim 1.5 \cdot 10^5$ between 970 and 1195 Å, and is in over-

all sensitivity limited to galactic studies (see e.g. Jenkins & Peimbert 1997).

Here we report on the detection at high spectral resolution of H₂ in the spectrum of the LMC star LH 10:3120. The star, located in the association LH 10 near the western edge of the LMC, is of spectral type O5.5Vf, has $V = 12.80 \text{ mag}$ and $E(B - V) = 0.17 \text{ mag}$ (Parker et al. 1992). The star was selected because of its very early spectral type, modest extinction (too large extinction would make the far-UV undetectable), and its proximity to an area where the molecule CO has been found in emission (Cohen et al. 1988; Israel et al. 1993).

The presence of H₂ in the LMC is known since Israel & Koornneef (1988) detected the near-IR emission lines from radiatively excited H₂ seen toward H II regions near hot stars. Measurements showed that H₂ is abundantly available, both in the SMC (Koornneef & Israel 1985) and in the LMC (Israel & Koornneef 1991a, 1991b). Clayton et al. (1996) detected with *HUT* at 3 Å resolution broad depressions in the far-UV spectrum of two LMC stars, which could be fitted with H₂ absorptions due to $N(\text{H}_2) \simeq 1\text{--}2 \cdot 10^{20} \text{ cm}^{-2}$. Studies of H₂ in the LMC and SMC are of importance because of the lower metal content of these galaxies compared to the Milky Way and the different gas to dust ratios (see Koornneef 1984).

2. Observations, data handling

The total observing time for LH 10:3120 in the *ORFEUS* space shuttle mission of Nov./Dec. 1996 was 6000 s in 3 pointings exploiting the integrating capabilities of the microchannel plate detector system. A detailed instrument description and information about the basic data reduction is given by Barnstedt et al. (1998). The data reduction for the 20 echelle orders has been performed by the *ORFEUS* team in Tübingen. The spectrum has been filtered by us with a de-noising algorithm basing on a wavelet transformation (Fligge & Solanki 1997). This leads to a slight degradation of the spectral resolution, now being equivalent to $\simeq 30 \text{ km s}^{-1}$. The spectra have a signal-to-noise (S/N) of $\simeq 20$ at the longer wavelengths of the recorded spectral range. Toward the shorter end of the spectral range, both the increased

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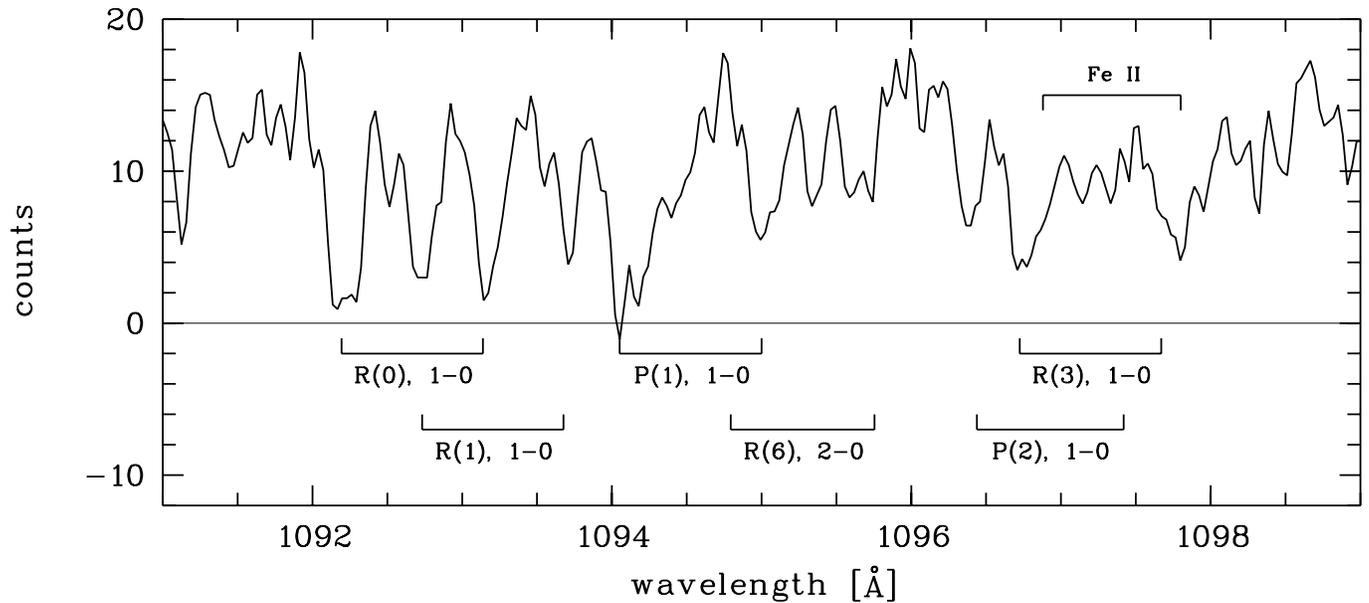


Fig. 1. A portion of the *ORFEUS* spectrum of LH 10:3120 near 1095 Å shows several H₂ absorption lines. The strongest of these are identified by their transition below the spectrum. The markings indicate the velocity range of 0 to 270 km s⁻¹ for each of the lines. Some weaker H₂ lines as well as some atomic features are also visible but have not been marked

Table 1. H₂ column densities in LMC gas toward LH 10:3120

Rotation level J	$\log N(J)$ [cm ⁻²]	b -value [km s ⁻¹]	Number of lines used
0	18.65	5	2
1	18.10	5	5
2	17.70	5	3
3	17.50	5	3
4	16.70	5	3

effect of the UV-extinction as well as the wavelength dependent sensitivity of the instrument leads to a reduction in S/N such, that little can be done with the spectrum at $\lambda < 1000$ Å. After the filtering absorption features in the longer part of the spectrum become clearly visible.

We first inspected spectral ranges almost devoid of atomic absorption lines and where only few H₂ absorptions are expected. The reason for this is to avoid getting confused in the search for H₂ by the complexity of the absorption line profiles on the line of sight to the LMC. Yet, the characteristic pattern of absorption by the Milky Way disk near 0 km s⁻¹, by the LMC near +270 km s⁻¹, and possibly by high-velocity clouds near +60 and +130 km s⁻¹ all known from *IUE* (see Savage & de Boer 1979, 1981; de Boer et al. 1980) and *HST* (Bomans et al. 1995) spectra, helps to identify the absorption structures.

A section of the echelle order 51, where several H₂ absorption lines have been found, is shown in Fig. 1. Many of the H₂ profiles overlap in their ~ 300 km s⁻¹ wide profile structure and decompositions are not always possible. However, in some cases the galactic absorption stayed unblended, in other cases the LMC portion was blend free. For this first analysis

we took 16 H₂ absorption lines from the lowest 5 rotational states for the further analysis. These lines are essentially free from any blending problems so that wrong identifications can be excluded. Absorption strengths could thus be determined for the LMC components seen in these absorption profiles. A selection of characteristic H₂ absorption line profiles in velocity scale (LSR) is shown in Fig. 2.

3. H₂ column density in the LMC gas

The absorption profiles have been analysed and decomposed into the various velocity components. Here we limit ourselves to the absorption pertaining to the LMC. For each line the absorption equivalent width has been determined. The f -values for the further analysis have been taken from Morton & Dinerstein (1976) for the H₂ transitions and for the atomic absorption lines from the compilation of Morton (1991). Theoretical Voigt profiles were fitted to some of the identified H₂ absorption lines and we thus could derive an upper limit for the velocity dispersion of $b < 10$ km s⁻¹ for the LMC gas.

Subsequently, curves of growth have been constructed for each of the absorptions by the 5 rotational states for the LMC component. The logarithmic equivalent widths for each rotational level J have been shifted horizontally to give a fit to a theoretical single cloud curve of growth as a function of $\log N(J)f\lambda$. The best fit for all 5 rotational states was obtained with $b = 5$ km s⁻¹ as shown in the empirical curve of growth for all identified H₂ lines (Fig. 3). The column densities $N(J)$ derived in this way have been collected in Table 1. The uncertainties in the column densities are based on the quality of the fits to the curve of growths. They range from 0.2 to 0.4 in the logarithm, as shown in Fig. 4. The total column density for the lowest 5 rotational states is $N(\text{H}_2)_{\text{total}} = 6.6 \times 10^{18}$ cm⁻².

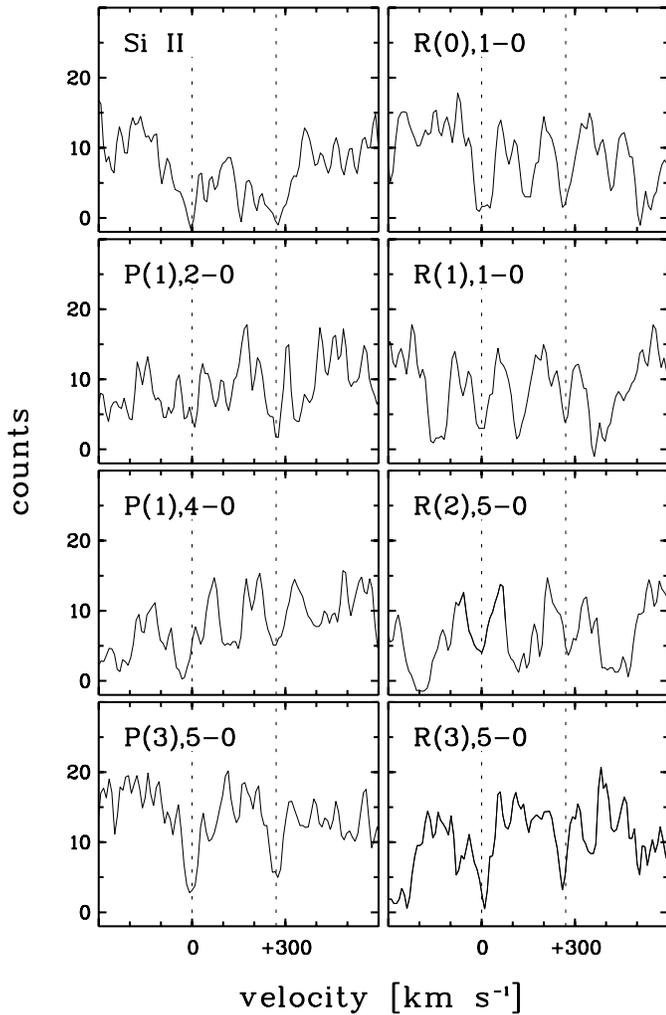


Fig. 2. Examples of absorption profiles clearly showing the detection of H₂. At the top we present a profile of the Si II 1190.416 Å line for easy reference. The H₂ absorption is due to the Milky Way component near 0 km s⁻¹ and the LMC component near +270 km s⁻¹. The atomic line of Si II also shows the presence of the galactic halo high-velocity clouds near +60 and +130 km s⁻¹ (LSR)

4. Interpretation

4.1. Excitation state

To explain the column densities and the relative population of the rotational states of the H₂ molecule we have to determine the mechanism which is responsible for the excitation of the molecular gas on our line of sight. Two processes dominate the population of the excited states, collisional excitation and pumping by UV photons (Spitzer & Zweibel 1974). With the column densities for the individual rotational states we were able to determine the excitation temperature for the LMC gas for $J \leq 4$. In Fig. 4 the LMC column densities $N(J)$ divided by their statistical weight g_J are plotted against the excitation energy E_J . The equivalent excitation temperature of the gas can be derived by fitting theoretical population densities from

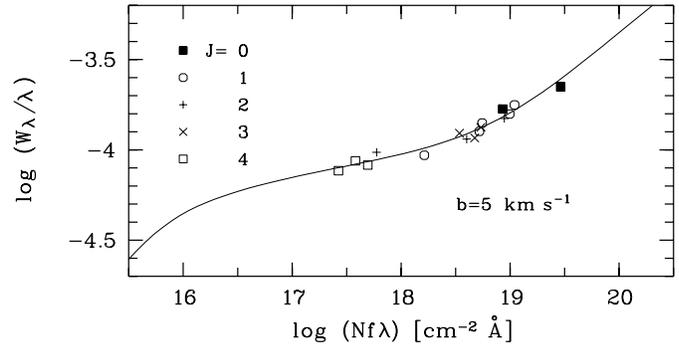


Fig. 3. The H₂ lines were fitted to a single cloud curve of growth, indicating $b \simeq 5$ km s⁻¹. The column densities for the levels $J = 0$ to 4 are given in Table 1

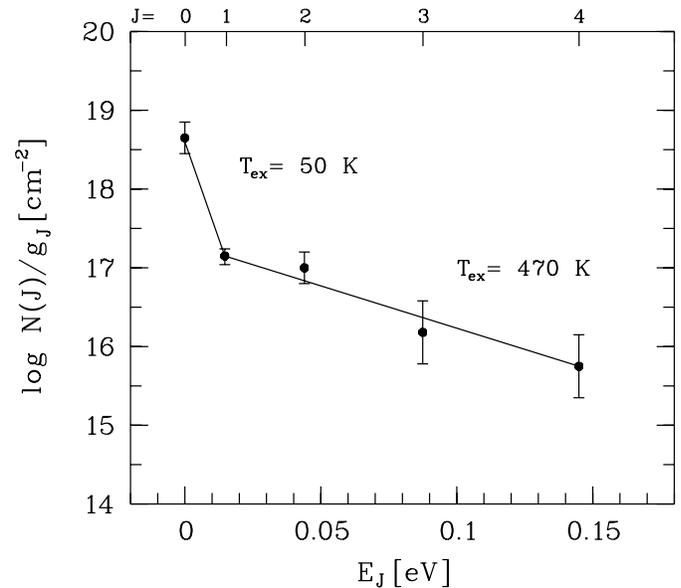


Fig. 4. The column densities for the levels $J = 0$ to 4 of H₂ are plotted against the excitation energy. The higher levels exhibit an equivalent excitation temperature of $\simeq 470$ K, the excitation being mostly due to UV pumping. For the two lowest levels the kinetic excitation temperature is ≤ 50 K; the fit drawn marks the indicated upper limit

a Boltzmann distribution. For $J \geq 2$ we derive an equivalent excitation temperature of 470 K.

The plot of Fig. 4 shows also that the level $J = 0$ lies above the linear relation of 470 K. Thus at least the $J = 0$ level is populated through kinetic excitation. Fitting the Boltzmann relation to the $J = 0$ and 1 column densities we find a kinetic temperature of 50 K. Clearly, the population of level $J = 0$ is due to gas with a temperature ≤ 50 K. Such very low gas temperatures have also been found by Dickey et al. (1994) and by Marx-Zimmer et al. (1998) in the LMC through analysis of H I 21 cm absorption.

Israel & Koornneef (1988) looked for emission by H₂ in the direction of N 11. They report a marginal detection of H₂ in the 1-0 $S(1)$ line at a position very close to our line of sight.

However, given the difference in direction as well as the low significance of the emission detection a connection with our absorption will be speculative. Also, the emission may be from gas more distant than the star.

4.2. Particularities of the line of sight

Our line of sight through the LMC is ending at the star LH 10:3120. The association LH 10 ionizes the H II region N 11B, the northern H II clump of the N 11 superbubble complex. The central superbubble is apparently created by the association LH 9 and is filled with hot gas. Several more patches of hot gas, partly coinciding with filamentary H α shells, exist. The one X-ray patch coinciding with N 11B hints at wind driven bubbles around some massive stars inside N 11B (Mac Low et al. 1998). While N 11B indeed shows first signs of the local effects of its most massive stars (Rosado et al. 1996), it is still a relatively unevolved H II region without large-scale expansion. Our line of sight is therefore most likely illuminated by the stars of LH 10. With the relatively unevolved nature of N 11B it is tempting to relate the H₂ absorbing gas with the remainder of the cold molecular cloud which formed N 11B.

4.3. Comparison with galactic gas

The total column density found in the LMC can be related to the properties of the H₂ gas in the Milky Way. The extinction to LH 10:3120 is $E(B - V) = 0.17$. Taking out the galactic foreground extinction of $E(B - V) = 0.05$ (Parker et al. 1992), one has extinction in the LMC of $E(B - V) = 0.12$. The total H₂ column density in the LMC is, compared to Milky Way gas (see Fig. 4 of Savage et al. 1977), slightly on the low side for its extinction. The smaller dust to gas content of the LMC (compared to the Milky Way; see Koornneef 1984) may have influenced the H₂ to $E(B - V)$ ratio by way of a lower H₂ formation rate.

The excitation level of the higher J levels on the line of sight to LH 10:3120 is equivalent to $\simeq 470$ K. Such values are also found for the higher rotational states in galactic gas (Snow 1977; Spitzer et al. 1974). The environment of N 11B has contributed to the UV pumping but not in an excessive way. The excitation of the lowest level indicates the gas is kinetically cold.

5. Concluding remarks

Our investigation of the *ORFEUS* far UV spectrum shows for the first time well resolved absorption profiles due to H₂ in LMC gas. The derived column densities and excitation temperatures toward LH 10:3120 demonstrate that UV pumping also takes place in LMC gas. The detected H₂ has a column density and an excitation state similar to that of Milky Way gas, even though the line of sight runs through an energetic environment. In closing we note that H₂ has also been detected in absorption by *ORFEUS* in the SMC (Richter et al. 1998).

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